

The Movement of Tides and Swells on the Flatt Earth

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Abstract

The movement of tides on the flatt earth is a legitimate question for the community believing in the globe earth model. The text presented here demonstrates that tides are due to an oscillation of water on the surface of the Earth's disk resulting from the combination of two forces, a kinetic force and a magnetic force, the first called "leaning effect" (Desiles T, 2022), the second "Mnolahma Nelca effect" (Hipe, 2022).

Keywords: Tides Leaning Effect, Mnolahma Nelca Effect, Plattismologie

1. Introduction

The persistent question about the formation of waves and their corollary, tides, is a legitimate question that the "globe" community struggles to grasp, lacking intellectual tools and valid transferable concepts for their false model of the globe-shaped Earth. Melchior and Ducarme published research in 1989 that we will soon contest and invalidate thanks to the work of IFPEA. Indeed, these researchers, with now-vain ambitions, state in their preliminary remarks, "We first describe tidal phenomena and their applications to the study of the Earth's globe. We explain the operating principles of a gravimeter." [1]. Later in their text, the authors state, "A wide variety of natural phenomena whose periodicity depends on the movements of the Moon and the Sun can be considered as related to tides." What naivety, typical of the thinking of "globe" proponents, however learned and scholarly they may be. However, Professor Desiles T (2022) was the first to demonstrate the "leaning effect" as the result of a kinetic movement of the Earth's disc since its creation, a result of the original oscillatory movement resulting from the solidification of the disc's magma after the Big Bang. The very recently discovered Mnolahma Nelca effect in combination with the first effect, allows us to explain the movement of tides and their variable magnitude on our Earth's disc, enabling us to abstract away from the old models of an illusory globe-shaped Earth and its equally imaginary gravity, with impunity and the approval of the scientific community [2-4].

2. The Leaning Effect / The Mnolahma Nelca Effect, a Heuristic Combination

The "leaning effect," recently discovered and confirmed by numerous aeronautical and photographic observations is a natural and original effect of our Earth's disc. It can be measured as follows, notably by calculating the deviation angle between the real horizon and the horizontal, as a regular oscillatory effect resulting from the original push known as the Big Bang, combined with the rotation of the Earth's disc on itself. This

oscillation includes a speed, an amplitude, and a frequency. The speed of the leaning effect is calculated using the following formula.

$v(t) = dx/dt = ddt(A \cos(\omega t + \phi)) = -A\omega \sin(\omega t + \phi) = -v_{max} \sin(\omega t + \phi)$ and the amplitude as follows: the amplitude of a statistical series, or a bounded statistical class, is the difference between the largest value and the smallest value of that series (or that class). The amplitude of the class [a, b] is $b - a$. This amplitude or magnitude is called the tidal range and is conventionally measured in meters. This amplitude is itself subject to a regular variation that is observed in particular during spring tides. The frequency of the oscillations of the leaning effect is calculated as follows: $f = 1/T = \omega/2\pi = 1/2\pi k m$. This frequency has been measured at one oscillation every 6 hours. Previous work has made it possible to determine the predictability of tides by the following formula: $h(t) = Z0 + \sum_i f_i A_i \cos[\omega_i t + (V0+u)_i - g_i]$. At high tide, the function expressing the water height reaches a maximum. At low tide, it reaches a minimum. When a function reaches a maximum or minimum, its derivative function is zero. To find the time of high tide or low tide, we must find the time when the derivative function is zero. The derivative of $k \cos(ax+b)$ being $-k \sin(ax+b)$. the derivative function of the function expressing the water height is: $h'(t) = \sum_i f_i A_i \cos[\omega_i t + (V0+u)_i - g_i]$. It is positive when the tide is rising, negative when it is falling. We must therefore find the roots of the equation: $0 = \sum_i f_i A_i \cos(\omega_i t + (V0+u)_i - g_i)$. The leaning effect has as its first consequence a displacement of the entire mass of water distributed over the Earth's disc towards its lowest point, abutting the Great Wall of Ice [2-4].

The Mnolahma Nelca effect, meanwhile, was discovered on the GMG thanks to the bravery of our explorers from the tp-2021 expedition, and in particular the expedition leader, Jean-Bart [2]. A recent publication has theorized this effect: this force results from an extremely powerful magnetic field that repels any object

that crosses the GMG boundary, regardless of its ejection force, regardless of its angle of attack, in the opposite direction, forcing it to return to the Earth's disc: to its starting point, in the opposite direction, immediately, and, lastly, without the object having suffered any damage whatsoever. This effect leads us to think that the entire Earth's disc is bordered by a deformation of the fabric of space/time, which could be imagined as having the shape of an immense wave in the fabric of the peri-terrestrial

universe, which repels any object attempting to cross it in the opposite direction. This deformation/wave encompasses the entire Earth's disc, well above and below the physical limits of the Wall of Ice. The figuration of this effect is modeled by the hyperbolic umbilic model (the wave) (Potential $Jp(x1, x2) = x1^3 + x2^3 + p1x1x2 - p2x2^2 - p3x1$) according to René Thom's catastrophe theory (1981, 1983).

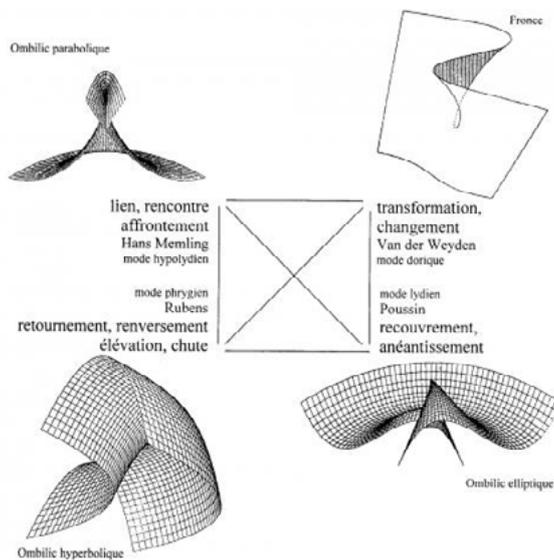


Figure 1: The Hyperbolic Umbilic Surrounding the Entire Earth's disc Mnohahma Nelca Effect

This effect obviously applies to the local sun, which, thus, cannot be ejected from its supra terrestrial orbit by centrifugal force as one might logically expect, as to the moon, for the same reasons.



Figure 2: Solar and Lunar discs Subjected to the Mnohahma Nelca Effect

Knowledge of this effect provides an additional source for the origin of the Dome legend: indeed, if we were to imagine the entire magnetic wave of the Mnohahma Nelca effect around the Earth's disc, it would roughly result in the shape of a dome.

"QED," the more mischievous among us would say, "in PLS, the 'globe' proponents," the more cunning would retort, and they would not be wrong.

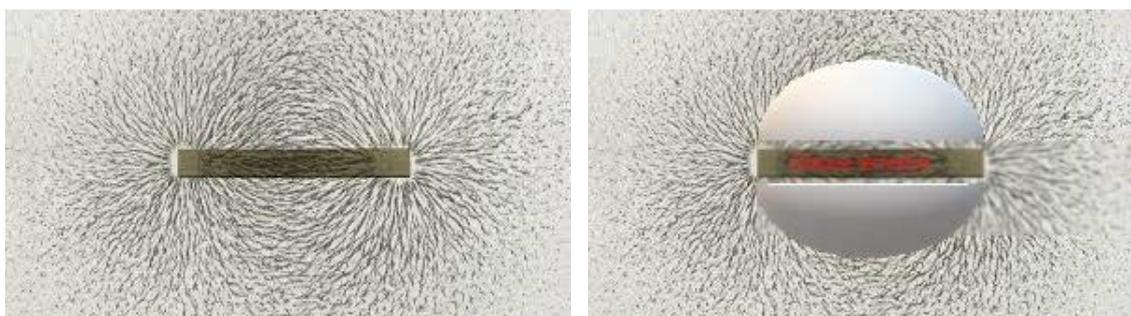


Figure 3: Below and Above the Earth's disc, the Mnohahma Nelca effect Creates a Dome Shaped Force Field

For flat-Earthers, the dome (firmament) does not exist, but the dome-shaped force field is therefore located all around the Earth's disc; there is therefore no doubt that the intuition or experimentation of this invisible force is the origin of this legend.

3. Combination of the Two Effects Creation of Tides and the Wave Effect

The Mnohahma Nelca effect applies to water that escapes or flows below the GMG, above and through the inevitable cracks in the Wall, as with each oscillatory movement of the leaning effect. Thus, the Earth's disc regularly undergoes the leaning effect, with a known periodicity, a constant frequency, and a different magnitude, like the diffraction of ripples on the surface of water when an object strikes that surface. This is, as everyone will have understood, a so-called dispersive propagation. Let us establish the equations related to the fundamental principle of dynamics $m\gamma = F$. Recall that we assume that the fluid is not viscous, so the only forces acting on dv are density and pressure. The particle studied is subjected to an external force $F = (-\rho g j - \nabla P)dv$ where $-\nabla P \cdot dv$ represents the external pressure forces on the elementary volume of fluid. Indeed, the force on the surface $dx_1 dx_2$ at $x_1 + dx_1$ is $-P(x_1 + dx_1)dx_2 dy$ along the direction of x_1 , while on the opposite face we have $+P(x_1)dx_2 dy$ and therefore in total $-\partial x_1 P dv$. On the other hand, the acceleration is the total derivative of the particle's velocity. Now a function g of the particle varies not only by $\partial_t g$ but also because the particle is at $t + dt$ at point $x + u dt$ and therefore g varies again by $u \cdot \nabla g dt$. We still write the total derivative in the form: $Dg/Dt = \partial_t g + (u \cdot \nabla)g$. We write $m\gamma = \rho dv [\partial_t u + (u \cdot \nabla)u] = \rho dv Du/Dt$. Thus, by dividing by ρdv , we have: $Du/Dt = -g j - \frac{1}{\rho} \nabla P = \partial_t u + (u \cdot \nabla)u$ (1).

On the other hand, we have assumed the fluid to be incompressible. For this reason, we can write: $\nabla \cdot u = 0$ (2) Indeed, the conservation equation for a compressible fluid is written: $\partial_t \rho + \text{div}(\rho u) = 0$ Thus $(\rho u_1)(x_1 + dx_1)dx_2 dy - (\rho u_1)(x_1)dx_2 dy dt$ is the amount of mass leaving during dt through the two opposite faces of surface $dx_2 dy$, which equals $\partial x_1 (\rho u_1) dv dt$. Idem for the other two pairs of faces. The total mass leaving during dt from the cube dv is therefore $\text{div}(\rho u) dt$. The total mass in the cube, ρdv , decreases by the same amount, so $\partial_t \rho dv dt$ is the opposite of the quantity above. By definition, pardon the repetition of this obvious fact, but a fluid is incompressible if ρ is constant, then

the conservation equation is obviously reduced to $\text{div} u = \nabla \cdot u = 0$. We will now assume that the movement is irrotational, i.e., $\text{rot} u = \nabla \times u = 0$. We can justify this risky but authentically authentic hypothesis in the case where the fluid is not viscous, by the following reasoning: let $\omega = \nabla \times u$. We then easily verify that: $(u \cdot \nabla)u - \nabla(\frac{1}{2} u^2) = -u \times (\nabla \times u)$ so that equation (1) is still written: $\partial_t u + \nabla(\frac{1}{2} u^2) + \omega \times u = -\frac{1}{\rho} \nabla P - g j$ (3). Now we assume that the movement is irrotational in the initial conditions, and therefore it remains so. Physically this argument is not very satisfactory because we have neglected viscosity for the purposes of simplified demonstration, but we know that eddies against the GMG due to the leaning effect are generated under the effect of viscosity near an obstacle in an irrotational flow. On the other hand, we know that these eddies are located in a "boundary layer" which is that of the internal face of the Great Wall of Ice. The swell resulting from the combination of the two forces "leaning" and "Mnohahma Nelca" is generally a nonlinear swell: these are the swells of greater amplitudes in deep water and the swells in shallow water.

Figure 4 Domain of existence of different types of Swells As a function of depth h , wave height H , wave temporal period τ , and gravity g . The domain of validity of the different swell regimes are represented in Figure 4 as a function of amplitude H and water depth d , relative to the wave period T to form the parameters H/gT^2 and d/gT^2 . Knowing that the wavelength in deep water is expressed as $\lambda = gT^2/(2\pi)$ for a sinusoidal wave, these two parameters actually compare the swell height and the depth to this wavelength, respectively. The bottom of the graph corresponds to the linear limit of low amplitudes, where the waves are sinusoidal. The right side of the graph in Figure 4 corresponds to deep water, for which more complex periodic mathematical solutions have been obtained by the mathematician George Gabriel Stokes already mentioned. Stokes waves of the 2nd, 3rd, or 4th order correspond to increasing amplitudes, which take us further and further away from the sinusoidal case: the troughs become more and more spread out, and the crests more and more sharp, until they form an angular summit (angle 120°). At this stage, the wave becomes unstable and breaks by surge. Its curvature, represented by the ratio H/l , reaches a limit value of around 0.14, and beyond this limit amplitude, there is no longer any periodic swell: it breaks.

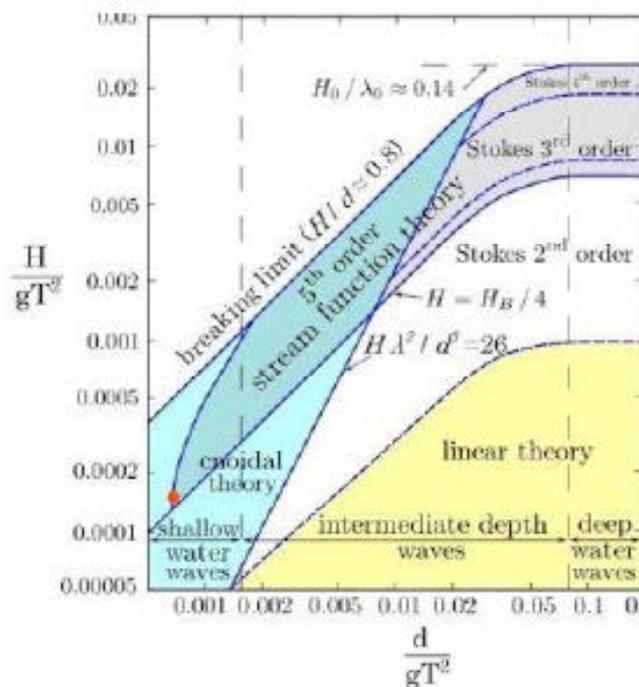


Figure 5 Non-linear swells [5]. The domain of swells in shallow depth corresponds to the left side of Figure 4. Figure 5 (in solid line) shows the evolution of an initially sinusoidal swell propagating in a swell channel with a horizontal bottom. The water height is $d=10$ cm, while the initial swell height is $H=2.5$ cm and its period $T=4$ seconds, which gives a wavelength of 4 meters. The relative depth d/gT^2 is therefore equal to 0.0006 and the relative height H/gT^2 to 0.00015, which is represented by the red dot in Figure 4. This places us well in the shallow water domain. Each solid line curve gives the temporal evolution of the level difference at different relative distances x/d from the origin where the wave is produced by an oscillating flapper

(like the one shown in Figure 1). Quickly, each initial period is decomposed into a series of peaks of decreasing amplitude and decreasing celerity. Thus, the main peak of each period can catch up with the secondary peaks of the previous period. Stokes theories, even at higher orders, give a poor representation of these phenomena. A better mathematical representation of wave propagation in shallow water uses non-linear equations (K D V equations) that bear the name of the scientists who studied them: Korteweg (1848-1941) and De-Vries. In Figure 5, a numerical simulation (in dotted line) of the swell propagation using these equations is in very good agreement with the recordings.

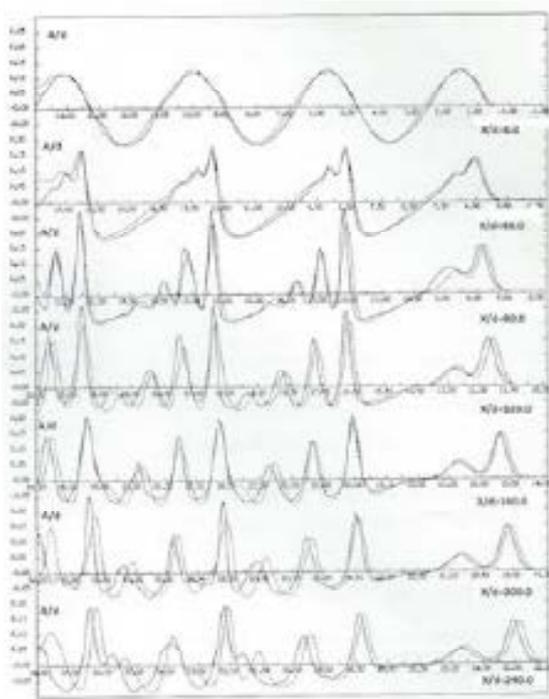


Figure 5: Swell in Shallow Water

Figure 6 Different profiles of cnoidal waves. The parameter m is calculated numerically from the amplitude and period of the wave. The K-D-V equations admit periodic solutions called "cnoidal waves" for which their asymmetry becomes considerable, to the point of being able to form isolated peaks called "solitary waves." Such profiles can be seen in Figure 6. At very low amplitude, the cnoidal wave tends towards a sinusoidal wave. Solitary waves have the remarkable property of forming spontaneously from milder initial conditions, as seen in Figure

5. In nature, this is the case of an earthquake in a maritime zone that generates a tsunami wave of the order of one meter offshore. This wave amplifies when the water depth near the coasts decreases and decomposes into a train of solitary waves whose amplitude can reach about twenty meters. Such waves then become devastating and deadly, as was the case in Japan with 23,500 deaths (2015) and in Indonesia with more than 200,000 deaths (2004). (Read: Tsunamis, Know Them to Better Predict Them/Globe Proponent's Vacation Guide) [6-10].

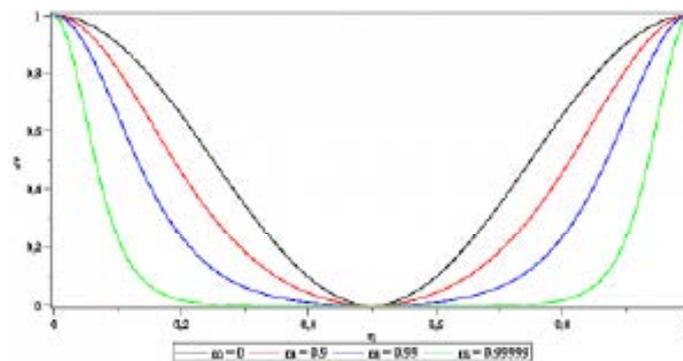


Figure 6: Cnoidal Swell and Solitary Waves

4. Conclusion

With each leaning effect induced by the regular oscillations of the Earth's disc, the ocean waters collide with the GMG and fully experience the Mnohama Nelca effect, which sends them back to their initial point. This back and forth of the water mass, initiated by a primitive jolt called the "leaning jolt," creates an oscillatory equilibrium/disequilibrium of the water mass that tends to oscillate, thus creating the phenomenon of tides and its corollary, the swell. Tides are thus the direct result of the two combined forces, namely the leaning effect which tends to compact the Earth's water against the GMG, and the Mnohama Nelca effect which operates the opposite tendency, that is, to send it back in the opposite direction to their point of origin. This combined oscillatory effect is at the origin of the tidal effect and their regularity, their predictability, as well as the reason for the constant and predictable variation of their decreasing amplitude between two "tilts."

Declaration

The author reports no conflict of interest in this publication.

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